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Photoreversible Birefringence Change of Diarylethene Single Crystals as Revealed by Change in Molecular Polarizability Anisotropy

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Abstract: Stimuli-responsive organic crystals represent a new frontier of material chemistry. Recently, we have reported photoreversible interference color change to multicolor in single crystals of а photochromic diarylethene derivative, 1,2-bis(2-ethyl-5-phenyl-3composed thienyl)perfluorocyclopentene (1a), accompanied by the photochromic reaction. The origin of the interference color change is due to the photoinduced birefringence change in the photoisomerization of diarylethenes. In this study, we newly found that single crystals composed of 1,2-bis(2,5-dimethyl-3-thienyl)perfluorocyclopentene (2a) also exhibit a photoreversible interference color change. The birefringence value for crystal 2a increased with the photocyclization conversion, while that for crystal 1a decreased. The relationship between the photoinduced birefringence changes for crystals 1a and 2a and their molecular structures was discussed based on the change in the molecular polarizability anisotropy accompanied by the photochromic reaction. These results would provide not only new opportunities for the application of photochromic crystals but also useful strategies for the design of crystalline materials that exhibit the desired birefringence change.

Introduction

Recently, organic crystals responding to external stimuli such as light, heat, electricity, gas, and force have been enthusiastically investigated because of their unique properties (e.g. mechanical effects,^{1–3} luminescent color change,^{4–8} self-healing,^{9–12} and gas absorption/release¹³). Photochromic diarylethene crystals have been investigated as the photoresponsive organic crystals having excellent photochromic properties such as high thermal stability of both isomers, high durability, high sensitivity, quick response, and high reactivity.¹⁴ Various physicochemical properties such as absorption spectra, refractive index, birefringence, permittivity, redox potential, fluorescence property, and electron spin in the crystal could be switched back and forth using UV and visible light.

Birefringence materials are essential to modulate the polarization of light. Thus, these are used in many application fields, such as the laser industry, optical communication, the polarimetry, and the scientific instrumentation.^{15–21} Organic crystals have birefringence due to molecular orientation. If birefringence in the crystal is switched by external stimuli, it should be expected as an innovative optical device like a switchable wavelength plate. Therefore, several groups have reported birefringence change in organic crystals by external stimulus. Horie et al. demonstrated that birefringence in the crystal composed of pseudorotaxane was reversibly changed by a phenyl ring rotation associated with a thermal phase transition.^{22,23} Moreover, Setaka et al. reported the birefringence change in crystalline molecular gyrotops.^{24–28} This is because a phenyl group encased in the three alkyl spokes was rotated. The higher the temperature of the crystal is, the faster the phenyl group rotates. Also, Diao et al. provided the interference color change accompanied by the birefringence change caused by a martensite phase transition of some organic semiconductor crystals.^{29–31} These birefringence modulations are induced by the change in temperatures such as heating and cooling processes. Moreover, Koshima and co-workers reported the photoreversible linear and circular birefringence change in the crystal of chiral salicylidenephenyletheylamines by employing the generalized high-accuracy universal polarimeter.³² In addition, very recently, we reported the photoreversible interference color change in single crystals composed of a diarylethene derivative, 1,2-bis(2-ethyl-5-phenyl-3-thienyl)perfluorocyclopentene (1a) (Scheme 1).³³ The interference color change was originated from the birefringence change associated with the photochromic reaction, which was quite sensitive to the reaction. The birefringence modulation by light irradiation has advantages over that by temperature (heat) because no direct physical contact is required, and it is possible to control spatially and temporally. On the other hand, in these previous works about the birefringence change in organic crystals, the relationship between the birefringence change and molecular structure was not quantitatively investigated. To design the crystals that show the photoinduced birefringence change, it is necessary to find out the relationship between the birefringence change and the molecular structural change in detail.

In this study, we newly investigated the change in the photoreversible interference color of single crystals composed of a diarylethene derivative, 1,2-bis(2,5-dimethyl-3-thienyl)perfluorocyclopentene (2a) (Scheme 1). We report the photoinduced birefringence change of crystal 2a associated with the photochromic reaction and the relationship between the change in birefringence and their molecular structures in comparison with crystal 1a. Our results may provide not only a new possibility for practical applications of photochromic crystals but also strategies for the design of molecular crystals that exhibit the desired birefringence change.



Scheme 1. Molecular structures of photochromic diarylethenes used in this work

Results and Discussion

Interference Color Change for Crystal 2a. Thin single crystals of 2a with a thickness from a few to some tens of micrometers were successfully prepared by a sublimation method (Supporting Information). The crystal shape of 2a is shown in Figure 1a. The crystal system and space group of crystal 2a are monoclinic and $P2_1/c$, respectively.³⁴ Crystal 2a consists of six faces that can be classified into two different types of faces. One is "Surface A" corresponding to (011), (011), (011), and (011) Miller planes. The other is "Surface B" corresponding to (100) and (100) Miller planes.³⁴ Figure 1a and b show the optical microphotographs of single crystals of 2a well-developed on each surface and their molecular packings viewed from individual faces. Upon alternating irradiation with UV (365 nm) and visible light, both crystal surfaces changed color by the photochromic reaction between the colorless open-ring isomer (2a) and the colored closed-ring isomer (2b).



Figure 1. Crystal shape and face indices of crystal **2a** and optical microphotographs of single crystals of **2a** viewed from (a) Surface A (thickness: 20.7 μ m) and (b) Surface B (thickness: 4.6 μ m), and their molecular packings viewed from the flat faces. The scale bar is 10 μ m.

When the single crystals of 2a were observed by rotating the sample stage at 45° from extinction position under crossed Nicols, the both crystals well-developed on Surfaces A and B exhibited interference colors due to birefringence originated from the refractive index anisotropy in the crystalline phase. To investigate the refractive index anisotropies of crystal 2a on both surfaces and ascertain the magnitude of the refractive index in two orthogonal directions, crystal 2a was observed using a sensitive-tint plate and a 1/4 wavelength plate as shown in Figures 2 and 3, which is the same method reported previously.³³ In the case of Surface A observed using the sensitive-tint plate under crossed Nicols, the color of the crystal at +45° and -45° from extinction position was yellow and green, respectively. When the 1/4 wavelength plate was used instead of the sensitive-tint plate, the color of the crystal at +45° and -45° was bluish-green and pale red, respectively. This means that the refractive index in the direction shown by n_1 is larger than that in the direction shown by n_2 in Figure 2n. On the other hand, when the sensitive-tint plate was used on Surface B, the color at +45° and -45° from extinction position was blue and orange, respectively. With the 1/4 wavelength plate, the color of the crystal at +45° became darker. This indicates that the refractive index in the direction shown by n_2 in Figure 3n. From these results, the retardation (*R*) of the crystals without any plate was estimated to be approximately 870 and 80 nm for Surfaces A (thickness: 5.1 µm) and B (thickness: 2.0 µm), respectively. The *R* values depend on the crystal thickness on the same surface. However, because the accurate Δn values have a wavelength dispersion, they can be determined by measuring the transmittance spectrum as shown later.



Figure 2. (a-1) Photographs under open Nicol, under crossed Nicols, and through a sensitive-tint plate and a 1/4 wavelength plate at +45°, 0°, and -45° from extinction position, (m) the direction of polarizer and analyzer and the direction of x' and z' for the sensitive-tint and 1/4 wavelength plates, (n) refractive index anisotropy, (o) corresponding crystal packing diagram in crystal **2a** on Surface A (thickness: 5.1 µm), and (p) Michel-Lévy color chart.



Figure 3. (a-1) Photographs under open Nicol, under crossed Nicols, and through a sensitive-tint plate and a 1/4 wavelength plate at +45°, 0°, and -45° from extinction position, (m) the direction of polarizer and analyzer and the direction of x' and z' for the sensitive-tint and 1/4 wavelength plates, (n) refractive index anisotropy, (o) corresponding crystal packing diagram in crystal **2a** on Surface B (thickness: 2.0 µm), and (p) Michel-Lévy color chart.

Next, the change in interference color of the crystals accompanying the photocyclization reaction upon irradiation with UV light was examined without any plate. Surface A of crystal 2a with a thickness of 3.4 µm was used for the experiment. Upon irradiation with UV light, the interference color changed from violet to blue, green, yellow, and orange. The actual color of the crystal observed under open Nicol hardly changed (i.e., the color changed from colorless to pale red), as shown in Figure 4b. The interference colors were thermally stable at room temperature even after UV

irradiation was turned off. These results indicate that birefringence could be modulated by the photochromic reaction and the interference color change is quite sensitive to the photochromic reaction. Moreover, on Surface B (thickness: 14.5 μ m), the interference color changed from purple to blue and sky blue as shown in Figure 4c. Figure 4d shows the interference color chart with respect to the change in retardation. From the correlation between the color chart and the interference color change, it was revealed that the retardation for **2a** on both Surfaces A and B increased upon UV irradiation.



Figure 4. (a) The direction of polarizer and analyzer, (b) change in crystal color (under open Nicol) and interference color (under crossed Nicols) according to photocyclization reaction of crystal **2a** on Surface A (5 μ m × 18 μ m, thickness: 3.4 μ m) by irradiation with 365 nm light (Irradiation power: 0.4 mW cm⁻²), (c) change in crystal color (under open Nicol) and interference color (under crossed Nicols) according to photocyclization reaction of crystal **2a** on Surface B (20 μ m × 20 μ m, thickness: 14.5 μ m) by irradiation with 365 nm light (Irradiation power: 1.1 mW cm⁻²), and (d) Michel-Lévy color chart.

To quantitatively estimate the value of retardation from the interference color, the transmittance spectra of the crystal were examined. Figure 5 shows the transmittance spectra of the crystals observed by rotating at 45° from the extinction position under crossed Nicols. Before photoirradiation, the transmittance spectra exhibited clear interference curves due to birefringence. Upon irradiation with UV light, the transmittance spectrum on both surfaces of crystal **2a** was shifted to longer

wavelength. The resultant transmittance spectrum was used for fitting to theoretical formulas shown in eqs 1 and 2. Eq 1 is given as an expression representing the light intensity (I) transmitted through the birefringent material under crossed Nicols.³⁵

$$I = I_0 sin^2 (\pi R / \lambda)$$

= $I_0 sin^2 (\pi d \Delta n / \lambda)$ (1)

where I_0 is the transmitted intensity without the crystal from halogen lamp under parallel Nicols, *R* is retardation, λ is the wavelength of the light, Δn is birefringence, and *d* is the thickness of the crystal. Note that it is necessary to consider the wavelength dispersion of birefringence. Eq 2 represents the wavelength dispersion of birefringence, which is known as a Sellmeier-type equation.³⁶

$$\Delta n = A + \frac{B}{\lambda^2 - C} \tag{2}$$

where *A*, *B*, and *C* are coefficients. The details to estimate the value of birefringence are depicted in Figures S1 and S2, and the results are shown in Tables S1 and S2.



Figure 5. Transmittance spectra for (a) Surface A (thickness: 50.4 μ m) and (b) Surface B (thickness: 29.1 μ m) of crystal **2a** observed by rotating at 45° from extinction position under crossed Nicols. The power of UVirradiation with 365 nm light was 4 mW cm⁻².

Figure 6 shows the initial *R* and Δn values obtained by fitting for surface A (thickness: 50.4 µm) and surface B (thickness: 29.1 µm) of crystal **2a**. The value of Δn on surface A was larger than that on surface B. This is because the tilt of the long and short axis of the diarylethene molecule against the observed crystal face was different between surface A and surface B. These findings are clarified by the molecular polarizability determined by quantum chemical calculations in the later section. After UV irradiation, the birefringence is changed by the photochromic reaction as described above. In order to ignore the effect of the linear dichroism, the *R* and Δn values were determined in the wavelength region ($\lambda = 650-900$ nm) where **2b** has no absorption.³⁷ The *R* and Δn values of crystal **2a** on both Surfaces A and B along with the photochromic reaction from **2a** to **2b** were estimated as shown in Figure 7. The *R* values increased on both surfaces. Thus, the Δn values increased on both surfaces. These are consistent with the results of the change in the interference color described above.

The reversibility of birefringence by the reversible photochromic reactions was examined. Figure 8a shows the reversible transmittance spectral change on Surface A of crystal **2a** upon alternating irradiation with UV and visible light. When the reversibility of the transmittance change was observed at 665 nm, the reversible birefringence change was confirmed more than 5 cycles without any damage, as shown in Figure 8b. Figures 8c and d also show the reversibility of the photoinduced birefringence change on Surface B. Such reversibility is one of the advantages for the use of diarylethene derivatives with high durability character.



Figure 6. Initial birefringence values on both surfaces of crystal 2a.



Figure 7. *R* and Δn values at each UV irradiation time determined by fitting for (a, c) Surface A (thickness: 50.4 µm) and (b, d) Surface B (thickness: 29.1 µm) of crystal **2a**. The power of UV irradiation with 365 nm light was 4 mW cm⁻².



Figure 8. Photoreversible change in transmittance spectrum on Surfaces (a) A (thickness: 29.3 μ m) and (c) B (thickness: 34.9 μ m) of crystal **2a** upon alternating irradiation with 365 nm light (power: 100 mW cm⁻²) for (a) 1 s and (c) 5 s and visible light, and (b, d) the repeatability for transmittance observed at 700 nm.

Relationship between Birefringence Change and Molecular Structure. The question is why the interference color could be modulated by the photochromic reaction. One may think that it is due to a change in the thickness of the crystal upon photoirradiation (i.e., Photomechanical effect of the crystals upon photoirradiation). However, in the present case, the crystal thickness hardly changed because the conversion ratio from the open-ring isomer to the closed-ring isomer is quite low.³³ Here, we assume that it is due to the change in the molecular polarizability anisotropy. It is well known that there is a relationship between polarizability and refractive index as Lorentz-Lorenz equation (eq 3).³⁸

$$\left(\frac{n^2 - 1}{n^2 + 2}\right) \left(\frac{M}{\rho}\right) = \frac{4}{3}\pi N\alpha \tag{3}$$

where *n* is a refractive index, ρ is density, *M* is molecular weight, *N* is Avogadro constant, and α is polarizability. Eq 3 can be rearranged to eq 4 that represents the relationship between birefringence and molecular polarizability anisotropy.³⁸

$$\Delta n = \frac{2}{9}\pi \frac{(n^2 + 2)^2}{n} \frac{\rho N}{M} \Delta \alpha \tag{4}$$

where $\Delta \alpha$ is the molecular polarizability anisotropy (i.e., the difference in the molecular polarizability between two axes in the molecule; $\Delta \alpha = |\alpha_1 - \alpha_2|$). Eq 4 shows that the Δn is directly proportional to the molecular polarizability anisotropy $\Delta \alpha$. To estimate $\Delta \alpha$, we calculated the molecular polarizability tensor (α) by a quantum chemical calculation. Note that the molecular structures of the photogenerated closed-ring forms **1b** and **2b** in the crystals **1a** and **2a** were used for the calculation, which could be obtained as disordered components in the X-ray crystallographic analysis as shown in Figure S3.^{34,39,40} In **1b**, only the reactive carbon and sulfur atoms appear as disordered components. In contrast, all atoms in **2b** appear as disordered components. When the conversion ratio from the open-ring form to the closed-ring form in the crystal becomes high, the phenyl group in **1b** may rotate differently from the dihedral angle between phenyl and thiophene rings in the open-ring form.⁴¹ However, in this study, since the conversion ratio in the crystal is low, we assumed that all atoms except the reactive carbon and sulfur atoms are the same as those of **1a**.

The molecular structures, the axes of the output file (x, y, and z), and the polarizability tensor in the output files of the calculations are summarized in Figure 9 and Table 1. The quantum chemical calculation provides the molecular polarizability tensors (α) that are not diagonalized as shown in eq 5, which is difficult to handle as it is. Here, the polarizability ellipsoid was introduced. Generally, refractive index anisotropy is presented as the refractive ellipsoid.⁴² The intersection of the index ellipsoid with the observed plane is an ellipse and the refractive index (n_1 and n_2) can be obtained from the lengths of the semi-axes of the ellipse. Moreover, molecules are regularly aligned in crystal, and the refractive index is expressed as the sum of the polarizability.⁴³ Accordingly, the molecular polarizability anisotropy (α_1 and α_2) can be estimated by depicting the polarizability ellipsoid similar to the refractive index ellipsoid.⁴⁴ To depict the polarizability ellipsoid, diagonalization of the polarizability tensors as shown in eq 6 is required. The eigenvectors (s, k, and t) and the eigenvalues (α_s, α_k , and α_i) indicate the ellipsoid axes (the principal axes) and the lengths of the semi-axes of the ellipsoid (principal values), respectively.

$$\boldsymbol{\alpha} = \begin{pmatrix} \alpha_{xx} & \alpha_{xy} & \alpha_{xz} \\ \alpha_{yx} & \alpha_{yy} & \alpha_{yz} \\ \alpha_{zx} & \alpha_{zy} & \alpha_{zz} \end{pmatrix}$$
(5)
$$(\alpha_{xy} = \alpha_{yx}, \alpha_{xz} = \alpha_{zx}, \alpha_{zy} = \alpha_{yz})$$

After diagonalization,

$$\boldsymbol{\alpha} = (\boldsymbol{s} \quad \boldsymbol{k} \quad \boldsymbol{t}) \begin{pmatrix} \alpha_s & 0 & 0\\ 0 & \alpha_k & 0\\ 0 & 0 & a_t \end{pmatrix} (\boldsymbol{s} \quad \boldsymbol{k} \quad \boldsymbol{t})^{-1}$$
(6)

The vectors and values determined in this work are summarized in Table 2. The relationship between the principal axes and their molecular structures are also shown in Figure 9. As can be seen

from the figure, the ellipsoid axes mostly correspond to the molecular axes. All of the principal values for the long axes of diarylethene were largest. In the case of diarylethene **1a**, when the molecule was isomerized from **1a** to **1b**, the polarizability in the long axis direction of the molecule (α_s) largely increased from 513.166 to 834.944, and the polarizability in the short axis direction of the molecular (α_k) and the thickness direction of the molecule (α_t) slightly increased from 256.573 to 293.284 and from 242.583 to 254.897, respectively. This is because the conjugation length was changed by photocyclization, which had the most influence on the long axis direction of the diarylethene. In the case of diarylethene **2a**, when the molecule was isomerized from **2a** to **2b**, the polarizability in the direction of the long and short axes of the molecule (α_s and α_k) increased from 261.369 to 299.572 and from 200.573 to 224.417, respectively, while the polarizability in the thickness direction of the molecule (α_i) decreased from 142.194 to 123.959. The increase of the polarizability, this is because the molecular spatial spread of the open-ring isomer was larger than those of the closed-ring isomer.



Figure 9. The molecular structures and the axes of the output file (x, y, and z) in the output files of the calculations for (a) 1a, (b) 1b, (c) 2a, (d) 2b. The axes (s, k, and t) indicated the principal axes of molecular polarizability ellipsoid.

	$lpha_{_{\!X\!X}}$	$lpha_{xy}$	$lpha_{_{yy}}$	$lpha_{_{\!X\!Z}}$	$lpha_{_{\!Y\!Z}}$	$lpha_{_{ZZ}}$
1 a	505.966	0.235	256.287	43.369	2.048	250.019
1b	834.334	-2.357	293.122	-18.627	2.652	255.669
2a	259.628	-1.478	200.573	14.178	1.133	143.935
2b	242.372	-32.038	281.583	1.243	-2.095	123.993

Table 1. Polarizability (a.u.) in the output files of the calculations.^a

^a 1 a.u. = $1.649 \times 10^{-41} \text{ C}^2 \text{ m}^2 \text{ J}^{-1}$

Table 2. The principal axes and the principal values (a.u.) of polarizability ellipsoid.^a

	$lpha_{ m s}$	S	$lpha_{ m k}$	k	$lpha_{ m t}$	t
		0.987		-0.025		-0.161
1a	513.166	0.002	256.573	0.990	242.583	-0.143
		0.163		0.141		0.977
		0.999		0.007		0.032
1b	834.944	-0.005	293.284	0.998	254.897	-0.067
		-0.032		0.067		0.997
		0.993		0.019		-0.120
2a	261.369	-0.022	200.573	1.000	142.194	-0.022
		0.120		0.025		0.993
		-0.489		-0.872		-0.007
2b	299.572	0.872	224.417	-0.489	123.959	0.012
		-0.014		-0.001		1.000

^a 1 a.u. = $1.649 \times 10^{-41} \text{ C}^2 \text{ m}^2 \text{ J}^{-1}$



Figure 10. Polarizability ellipsoid of (a) 1a, (b) 1b, (c) 2a, and (d) 2b.

Moreover, the molecular polarizability ellipsoids were depicted in Figure 10. The intersection of the index ellipsoid with the plane viewed from the observed crystal surface was determined to estimate the polarizability anisotropy (α_1 and α_2). The plane was calculated by rotating the polarizability ellipsoid by an arbitrary angle using a rotation matrix. Figure 11 shows the molecular structures for the open-ring isomers (**1a** and **2a**) and the closed-ring isomers (**1b** and **2b**), the axes of the molecular polarizability, and the ellipse considering the observed crystal surface. As the determined axes of the molecular polarizability were almost matched with the experimental optical axes as shown in Figures 2 and 3, the values of α_1 and α_2 in each molecule were determined from the intersection of the ellipse with the experimental optical axes and summarized in Table 3. The molecular polarizabilities in **1a** and **1b** are larger than those in **2a** and **2b**, respectively. This is due to the π -conjugation by a phenyl group. From the molecular polarizability ellipsoid, we can also explain the dependence of the initial birefringence values in crystal **2a** on Surface A and B. As shown in Table **3**, the $\Delta \alpha$ value for **2a** on Surface B is 44, which is smaller than the value on Surface A ($\Delta \alpha =$ 96). This is consistent with the experimental result that Δn on Surface A is larger than that on Surface B. Depending on the direction of the observation, the ellipse shapes and the polarizability values were

different. Furthermore, the molecular polarizability should change in proportional to the conversion ratio from the open-ring isomer to the closed-ring isomer. Then, the relationship between the molecular polarizability and the conversion ratio from the open-ring isomer to the closed-ring isomer was depicted as shown in Figure 12. The change in width between two straight lines directly means the difference in $\Delta \alpha$ by the photochromic reaction. Note that the conversion ratio from the open-ring form to the closed-ring form is quite low (less than 10%). In the case of diarylethene **1a**, $\Delta \alpha$ decreases with the conversion ratio from **1a** to **1b**, while, for diarylethene **2a**, $\Delta \alpha$ increases with the conversion ratio from **2a** to **2b**. These results mean that the Δn for **1a** and **2a** decreases and increases respectively, which is consistent with the results we observed. Thus, the relationship between the photoinduced birefringence change of crystals **1a** and **2a** and their molecular structures can be successfully explained based on the molecular polarizability anisotropy.



Figure 11. The molecular structures for the open-ring isomers and the closed-ring isomers, the axes of the molecular polarizability, and the ellipse considering the observation on the crystal surface for (a) **1a**, (b) **1b**, (c) **2a** (Surface A), (d) **2b** (Surface A), (e) **2a** (Surface B), and (f) **2b** (Surface B).

	α_1	α_2	$\Delta \alpha$
1a	298	256	42
1b	296	293	3
2a (Surface A)	251	155	96
2b (Surface A)	294	142	152
2a (Surface B)	157	201	44
2b (Surface B)	134	225	91

Table 3. Calculated polarizability (a.u.) in crystal packing.^a

^a 1 a.u. = $1.649 \times 10^{-41} \text{ C}^2 \text{ m}^2 \text{ J}^{-1}, \Delta \alpha = |\alpha_1 - \alpha_2|.$



Figure 12. Relationship between the molecular polarizability and the conversion ratio from the openring isomer to the closed-ring isomer for (a) 1, (b) 2 on Surface A, and (c) 2 on Surface B.

Conclusions

We have prepared the thin single crystals of **2a** and investigated the photoinduced birefringence change by the photochromic reaction. The crystals exhibited photoreversible interference color change originated from birefringence change by the photochromic reaction, which was quite sensitive to the photochromic reaction. The relationship between the birefringence change and molecular structure was revealed in the viewpoint from the molecular polarizability anisotropy estimated by quantum chemical calculations. These results provide not only new opportunities for the application of photochromic crystals but also useful strategies for the design of molecular crystalline materials exhibiting the desired birefringence change.

Experimental Section

General. Optical microphotographs of crystals were also observed using a Keyence VHX-500 digital microscope and UV irradiation was carried out using a Keyence UV-LED UV-400 (365 nm light). Interference color changes were observed using a Nikon ECLIPSE E600POL polarizing optical microscope, equipped with a video camera system. Crystal thickness was determined using a KEYENCE VK-8700 laser scanning microscope. UV light irradiation to the crystals was carried out using a super high-pressure mercury lamp (100 W; UV-1A filter (365 nm light excitation)) attached to the polarizing optical microscope. Visible light irradiation was carried out using a halogen lamp (100 W). Transmittance spectra were measured using the polarizing optical microscope connected with an Ocean Optics USB4000 fiber optic spectrometer. Quantum chemical calculation was performed using Gaussian 09 Rev. C.01 program package.⁴⁵ The molecular polarizability was calculated at DFT B3LYP/6-31G(d) level of theory using the coordinates obtained from the X-ray crystal structure analysis previously reported.^{34,39,40} The CCDC numbers of **1a**, **1b**, **2a**, and **2b** is 159602, 629166, 147109, and 147108, respectively.

Materials. 1,2-Bis(2,5-dimethyl-3-thienyl)perfluorocyclopentene (**2a**) was synthesized according to the procedures described in the literature.³⁷ The micrometer-size single crystals were prepared by sublimation of the compounds on the glass plates.

Observation of crystals. Optical photographs of crystals were observed with a digital microscope under non-polarized light. To investigate the refractive index anisotropy, the observation was carried out using the sensitive-tint and 1/4 wavelength ($\lambda = 546$ nm) plates with a polarizing optical microscope under crossed Nicols with polarizer and analyzer being orthogonal. The interference color change accompanying the birefringence change was observed under crossed Nicols. At the same time, change in color tone of the crystal by the photochromic reaction was observed under open Nicol with the removal of the analyzer.

Associated content

Supporting Information.

Preparation of thin crystal **2a**, Determination of the value of birefringence change in crystal **2a**, molecular structure used for the quantum chemical calculation, including Tables S1 and S2 and Figures S1-S3 (PDF).

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Notes

The authors declare no competing financial interest.

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